

Finite difference approximations to derivatives

(using Taylor's series expansions)

$$\boxed{\frac{\partial u}{\partial t}} = \boxed{D \frac{\partial^2 u}{\partial x^2}} \quad \text{1D transient heat equation}$$

Temporal derivative:
$$u_i^{n+1} = u_i^n + \Delta t \cdot \left(\frac{\partial u}{\partial t} \right)_i^n + O(\Delta t^2)$$

rearrange:
$$\left(\frac{\partial u}{\partial t} \right)_i^n = \frac{u_i^{n+1} - u_i^n}{\Delta t} + O(\Delta t)$$

Spatial derivative:

$$u_{i+1}^n = u_i^n + \Delta x \cdot \left(\frac{\partial u}{\partial x} \right)_i^n + \frac{1}{2} \Delta x^2 \cdot \left(\frac{\partial^2 u}{\partial x^2} \right)_i^n + \frac{1}{6} \Delta x^3 \cdot \left(\frac{\partial^3 u}{\partial x^3} \right)_i^n + O(\Delta x^4)$$

$$u_{i-1}^n = u_i^n - \Delta x \cdot \left(\frac{\partial u}{\partial x} \right)_i^n + \frac{1}{2} \Delta x^2 \cdot \left(\frac{\partial^2 u}{\partial x^2} \right)_i^n - \frac{1}{6} \Delta x^3 \cdot \left(\frac{\partial^3 u}{\partial x^3} \right)_i^n + O(\Delta x^4)$$

add and rearrange:
$$\left(\frac{\partial^2 u}{\partial x^2} \right)_i^n = \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{\Delta x^2} + O(\Delta x^2)$$

Finite differences: stability

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2} \quad \text{1D transient heat equation}$$

$$u_i^{n+1} = u_i^n + D \frac{\Delta t}{\Delta x^2} (u_{i+1}^n - 2u_i^n + u_{i-1}^n) + O(\Delta t^2, \Delta x^2)$$

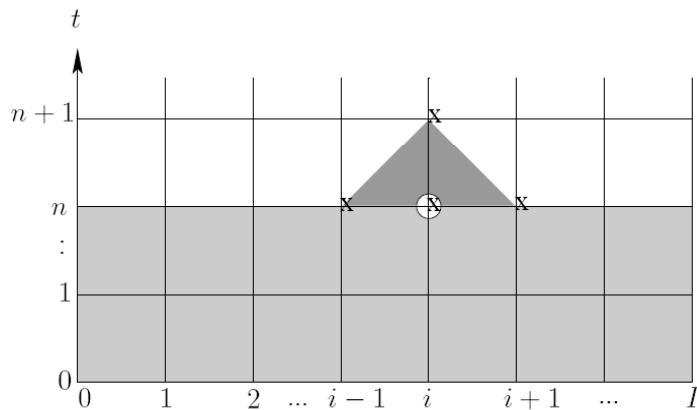
Conditionally stable when:

$$D \frac{\Delta t}{\Delta x^2} \leq \frac{1}{2} \quad (\text{Courant condition - Von-Neumann stability analysis})$$

Finite differences

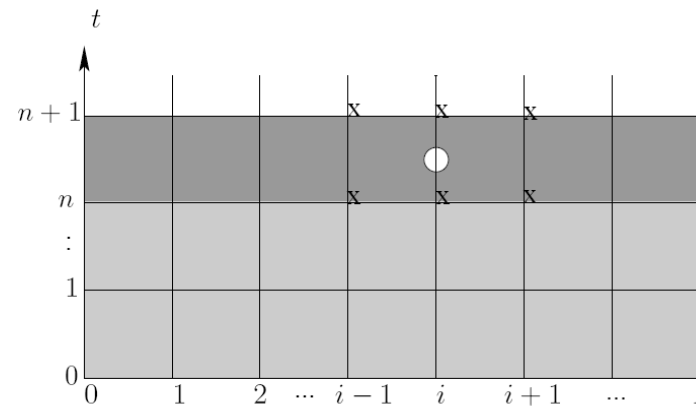
$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2} \quad \text{1D transient heat equation}$$

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = D \left\{ \theta \left(\frac{\partial^2 u}{\partial x^2} \right)_i^{n+1} + (1 - \theta) \left(\frac{\partial^2 u}{\partial x^2} \right)_i^n \right\}$$



Explicit forward differences

$$\theta = 0$$



Implicit central differences

$$\theta = \frac{1}{2} \quad \text{Crank-Nicolson}$$

Transient Advection-Diffusion FEM

$$\frac{\partial u}{\partial t} + \mathbf{v} \cdot \nabla u = D \nabla^2 u + f$$

Weighted residuals: $\int_{\Omega} \left(\frac{\partial u}{\partial t} + \mathbf{v} \cdot \nabla u - D \nabla^2 u - f \right) \omega \, d\Omega = 0$

Green-Gauss theorem (integ. by parts):

$$\int_{\Omega} \left[\left(\frac{\partial u}{\partial t} + \mathbf{v} \cdot \nabla u \right) \omega + D \nabla u \cdot \nabla \omega \right] d\Omega = \int_{\Omega} f \omega \, d\Omega + D \int_{\partial\Omega} \frac{\partial u}{\partial n} \omega \, d\Gamma$$

$$M \frac{du}{dt} + K u = F$$

FE approximation: split into elements and apply $u = \varphi_n u_n$ and $\omega = \varphi_m$

$$M_{mne} = \int_0^1 \varphi_m \varphi_n J \, d\xi \quad K_{mne} = \int_0^1 D \frac{\partial \varphi_m}{\partial \xi_i} \frac{\partial \varphi_n}{\partial \xi_j} \cdot \frac{\partial \xi_i}{\partial x_k} \frac{\partial \xi_j}{\partial x_k} J \, d\xi + \int_0^1 v_j \varphi_m \frac{\partial \varphi_n}{\partial \xi_i} \frac{\partial \xi_i}{\partial x_j} J \, d\xi$$

Transient Advection-Diffusion FEM

$$M \frac{du}{dt} + K u = F$$

Temporal integration: finite differences

Spatial integration: finite elements

Finite difference approximation to time derivative (Crank-Nicolson):

$$M \frac{u^{n+1} - u^n}{\Delta t} + K [\theta u^{n+1} + (1 - \theta) u^n] = F$$

c.f. 1D finite differences: $\frac{u_i^{n+1} - u_i^n}{\Delta t} = D \left\{ \theta \left(\frac{\partial^2 u}{\partial x^2} \right)_i^{n+1} + (1 - \theta) \left(\frac{\partial^2 u}{\partial x^2} \right)_i^n \right\}$

Iterate through time, solving spatial FEM problem (system of equations) at each time step:

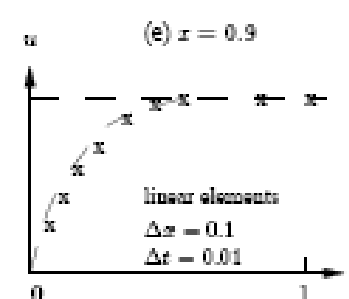
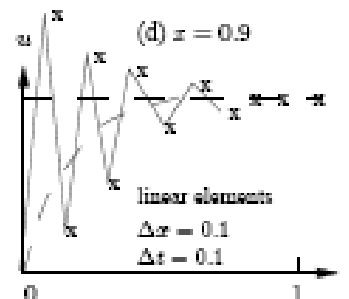
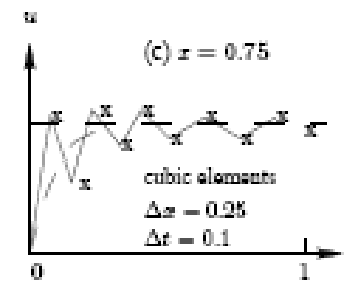
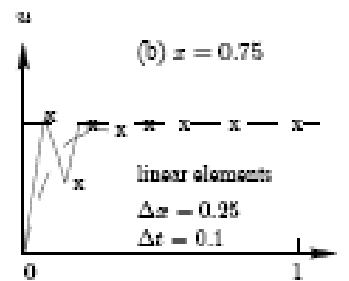
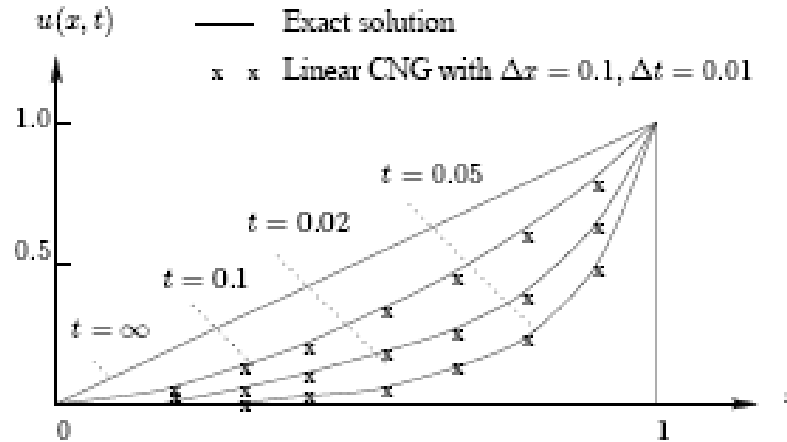
$$[M + \theta \Delta t K] u^{n+1} = [M - (1 - \theta) \Delta t K] u^n + \Delta t F$$

Transient Advection-Diffusion FEM: stability

Conditionally stable:

$$\Delta t < \frac{1}{(1 - \theta) \lambda_{\max}}$$

max eigenvalue of $M^{-1}K$.



Mass lumping

$$M \frac{du}{dt} + K u = F$$

Coupled system of ODEs, because **M** is fully populated.

So for explicit finite differences ($\theta = 0$):

we still need to solve a system of equations at every time step (slow!).

To decouple equations, “lump” row sums of **M** onto diagonals:

$$M = \begin{bmatrix} \frac{1}{9} & \frac{1}{18} & \frac{1}{18} & \frac{1}{36} \\ \frac{1}{18} & \frac{1}{9} & \frac{1}{36} & \frac{1}{18} \\ \frac{1}{18} & \frac{1}{36} & \frac{1}{9} & \frac{1}{18} \\ \frac{1}{36} & \frac{1}{18} & \frac{1}{18} & \frac{1}{9} \end{bmatrix} \xrightarrow{\text{mass lumping}} \begin{bmatrix} \frac{1}{4} & 0 & 0 & 0 \\ 0 & \frac{1}{4} & 0 & 0 \\ 0 & 0 & \frac{1}{4} & 0 \\ 0 & 0 & 0 & \frac{1}{4} \end{bmatrix}$$

Mass lumping + explicit finite differences ($\theta = 0$):

NO need to solve system of equations at every time step (fast!).

Mass lumping: phase lag errors

